

①.

Introduction

- ①. Will talk about 2D Yang Mills theory for $SU(N) / U(N)$ gauge group on Σ_G .

$$S = \int_{\Sigma_G} d^2\sigma \sqrt{g} (F_{\mu\nu} F^{\mu\nu})$$

This theory is invariant

under area-preserving diffeomorphisms

~~z~~

$$Z = \int DA e^{-S}$$

only depends on area & genus of $\Sigma_G(A)$.

$$f = *F \quad \Leftrightarrow \quad F = f_{\mu}$$

$$F_{ij} = \sqrt{\det G_{ij}} \epsilon_{ij} f^a$$

$$\int F \wedge *F = \int \mu f^2$$

$$\int_{\Sigma_G} \text{tr}(F \wedge *F) = \int_{\Sigma_G} \mu \text{tr}(f^2)$$

$$S_{YM} = \frac{1}{4g^2} \int_{\Sigma_G} d^2x \sqrt{\det G_{ij}} \text{Tr}(f^2)$$

$$= \left(\frac{1}{4g^2} \right) \int_{\Sigma_G} \mu \text{Tr}(f^2)$$

- $\text{Tr} f^2$

is a scalar

- invariant under diffeos

- As long as $\int \mu$ is invariant,

the action is invariant.

(2)

- gauge field in 2D: No local d.o.f.

- The exact solution is known:

For partition function

$$Z(G, A)$$

$$= \sum_R (\text{Dim } R)^{2-2G} e^{-g^2 A \Omega(R)}$$

Will write ans as

$$= \sum_R (\text{Dim } R)^{2-2G} e^{-A \Omega_2(R)}$$

Redefined: $g^2 A \rightarrow A$.

(3).

- Gross & Taylor (1992-1993) + others showed that the $\left(\frac{1}{N}\right)$

expansion of

$$Z(G, A)$$

$$= \sum_h \frac{1}{h!} N^{2-2h} Z_{\text{string}}(\Sigma_h \rightarrow \Sigma_G)$$

- $Z_{\text{string}}(\Sigma_h \rightarrow \Sigma_G)$

is a counting function of holomorphic

Maps

$$\Sigma_h \rightarrow \Sigma_G(A)$$

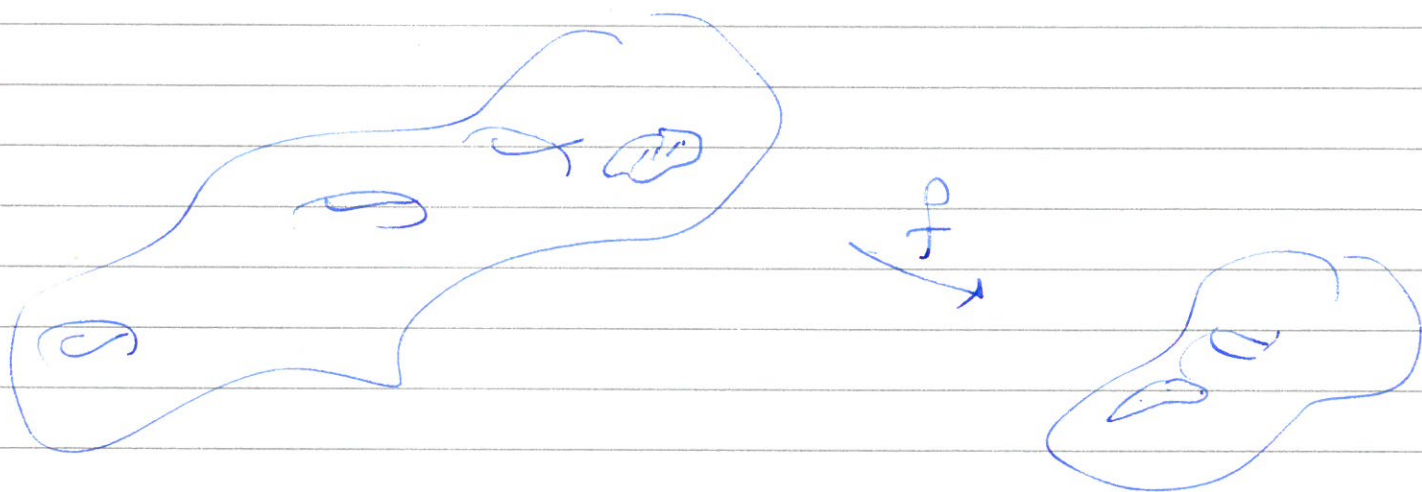
(4)

→ How ?

— what is a holomorphic
map from

$$\Sigma_h \rightarrow \Sigma_g \quad ?$$

— what does it look like ?



Local patch

$$z, \bar{z} :$$

$$f(z, \bar{z})$$

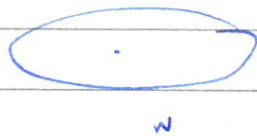
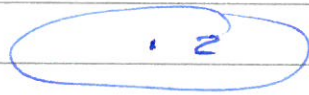
$$\partial_{\bar{z}} f = 0$$

(5)

Introducing local coordinate

z at P

& w at $f(P)$



$\mapsto w = z$ is a holo. map.

$w = z^n$ for $n \in \{1, 2, \dots\}$

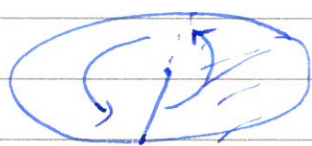
is a holo. map.

$\mapsto w = \sqrt{z}$, $\log z$ Not a hol. map.

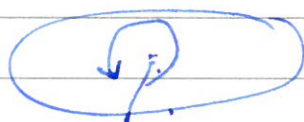
- Single-valued.

locally $(P, f(P))$

we have \equiv

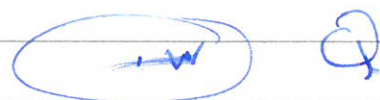


z^2



2-fold covers:

Fixing a point Q :



degree:

$k_1 + k_2 + \dots + k_n = N$

= Degree is constant as we move across Σ_g

7

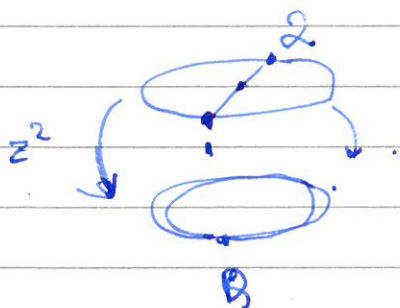
- Maps of a degree n
are classified by partition
of n :

~~is~~

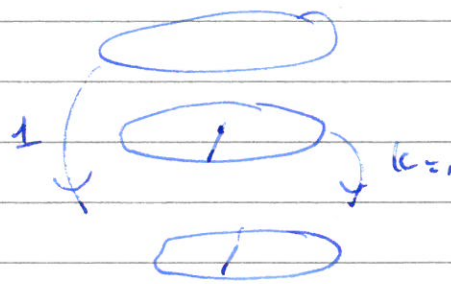
$$\begin{aligned} 3 &= 3 \\ &= 2+1 \\ &= 1+1+1 \end{aligned}$$

- Locally:

$$n = 2^2.$$



$$\begin{aligned} 1 &\rightarrow 2 \\ 2 &\rightarrow 1 \end{aligned}$$



$$\begin{aligned} 1 &\rightarrow 2 \\ 2 &\rightarrow 2 \end{aligned}$$

(8).

In general for degree n

maps: σ to a dotc,

we can construct a permutation

$$\sigma: \{1, \dots, n\} \rightarrow \{1, \dots, n\}$$

(3):

$$\begin{array}{ccc} 1 & 2 & 3 \\ \downarrow & & \\ 2 & 3 & 1 \end{array} = (123)$$

$$\begin{array}{ccc} 1 & 2 & 3 \\ \downarrow & & \\ 2 & 1 & 3 \end{array} = (12)(3) = (12)$$

$$\begin{array}{ccc} 1 & 2 & 3 \\ \downarrow & & \\ 1 & 2 & 3 \end{array} = (1)(2)(3) = ()$$

(9)

→ $(123) \neq (132)$

- describe the same map.

- choice of how you label
the points

$$\rightarrow (123) = \underbrace{(23)}_{\gamma} (123) \underbrace{(23)}_{\gamma^{-1}}$$
$$\sigma = \gamma \sigma \gamma^{-1}$$

→ locally looks like maps correspond
to perms.

→ Conjugacy classes of perms.

globally : Laks

Outline:

①. — Exact answer. (Migdal)

— Using heat kernel action

②. — How we go from $U(N)$

Representation sums to

Sums involving permutations —

which count Weyl morphisms

maps

— $\mathbb{C}[S_n]$ — group algebra.

③. — Centre of the group algebra.

— Schur-Weyl duality relation

between $U(N)$, S_n .

— Implications of S.W. for

③

• $\dim R$

• $\chi_2(R)$

③ Large \mathbb{N} Expansion of 2dYM
(- Chiral Exp.)

- Writing it in a form app. for
interp.

④ - Precise Mathematical Statement
of the conn. between forms of
holo maps

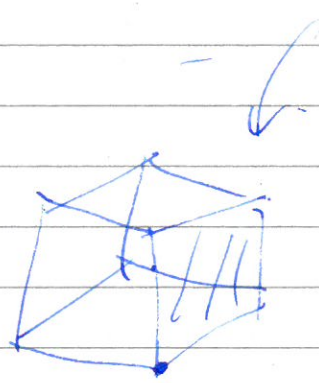
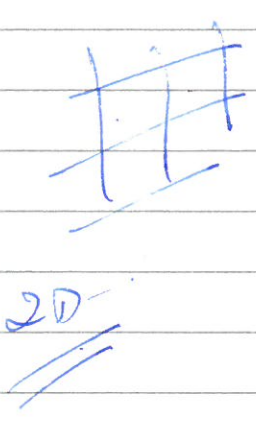
- How the large \mathbb{N} expansion
formulae can be seen to
be country maps,

Part I: Deriving the exact answer
for $Z_G(A)$:

⊗ Lattice gauge theory:

$$\int \prod_E dU_E \quad \prod_P \pi Z(U_P)$$

- Cubic lattice:



- plaquette.
- 2-cell.

4D:

x_1, x_2, x_3, x_4

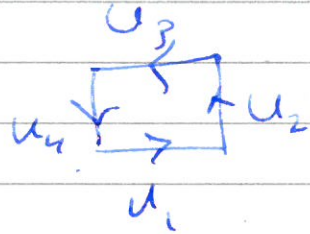
- (x_1, x_2)
- $(1, x_3)$
- (x_1, x_4)
- (x_2, x_3)

(i, j) : $i \in j$

Juraal

$$\rightarrow Z(U_p) = e^{-\frac{g^2 \text{Tr}(2 - U_p - U_p^\dagger)}{2g^2}}$$

$$U_p = (U_1 U_2 U_3 U_4)$$



→ Wilson action:

→ one shows: $\int_{\mathcal{G}} a \rightarrow 0$ (naive continuum limit)
 • discretizes T.N. $\int_{\mathcal{G}} \text{Tr} F^2$

Also: heat kernel action

- Shown to be equivalent to

Wilson at weak coupling

$$g \rightarrow 0$$

(13)

$$Z_H(U, A) = \sum_R e^{-g^2 A C_2(R)} \text{Dim } R \chi_R(U)$$

- Sum is over all irreps. of

The gauge group

- Dim R . - Dimension of the irrep

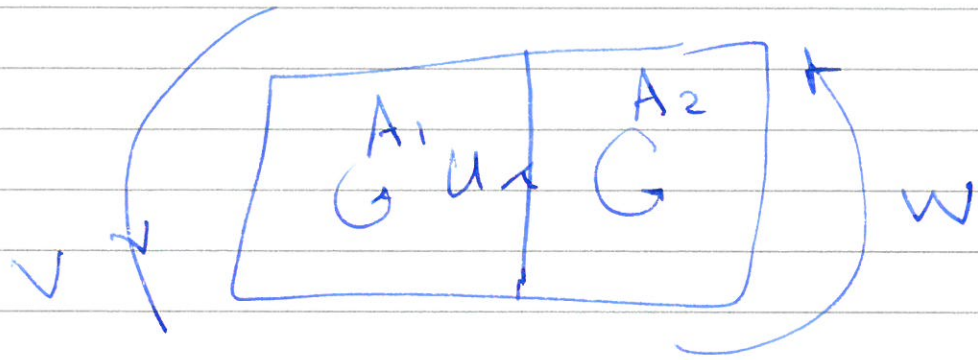
- $C_2(R)$ - Quadratic Casimir

- $\chi_R(U)$:

Character of U

$$\text{Tr}(\rho^R(U))$$

• This lattice partition function \rightarrow
 R.G. invariant



$$\sum_{R_1, R_2} \left(\chi_{R_1}(V u^+) \chi_{R_2}(u^+ W) \right) e^{-A_1 C_2(R_1)} e^{-A_2 C_2(R_2)}$$

$$= \sum_{R_1, R_2} \delta_{R_1, R_2} \frac{\chi_{R_1}(V W)}{\text{Dim } R_1}$$

$$\text{Dim } R_1 \text{ Dim } R_2 e^{-A_1 C_2(R_1)} e^{-A_2 C_2(R_2)}$$

$$= \sum_R (\text{Dim } R) e^{-A C_2(R)} \chi_R(V W)$$

$$A = (A_1 + A_2)$$

(I.5)

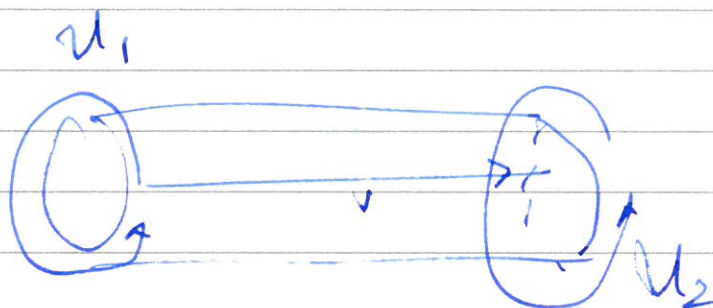
~~$Z_p(A, A)$~~

$$\begin{aligned} \mapsto \int & Z_p(A_1, u) Z_p(A_2, u^+ w) \\ &= Z_p(A_1 + A_2, v w) \end{aligned}$$

∴ We can use the coarsest
possible discretization —

— Will agree with a fine discretization
— i.e. with continuum limit.

Cylinder:



$$\sum_R \int dV \chi_R(u_1, v, u_2, v^+) \left(\text{Dim } R \right) e^{-A Q_2(R)}$$

$$= \sum_R \frac{\chi_R(u_1) \chi_R(u_2) \text{Dim } R}{\left(\text{Dim } R \right) e^{-A Q_2(R)}}$$

$$= \sum_R \chi_R(u_1) \chi_R(u_2) e^{-A Q_2(R)}$$

⊗ ~~the~~ fall

→ Haar Measure

→ G-invariant group
integ.

→ normalized to 1.

⊗.
$$\left\{ \begin{array}{l} \int d\mathbf{u} U_{ij} U_{kl}^{\dagger} = \delta_{ik} \delta_{jl} \\ \int d\mathbf{u} \delta_{ij} \delta_{kl} \\ \int d\mathbf{u} = 1. \end{array} \right.$$

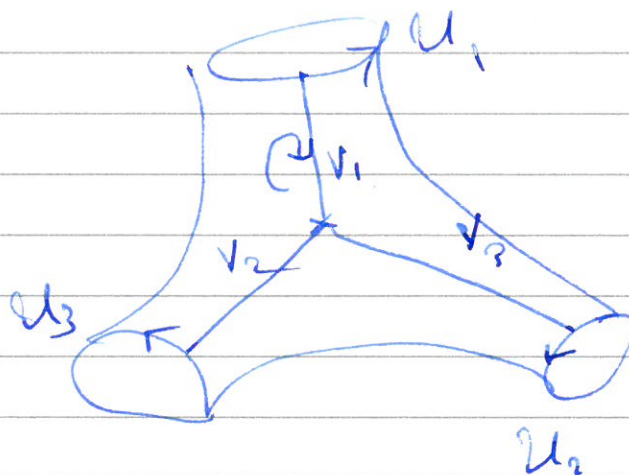
⊗.
$$Z(T^2) = \int Z_{\text{cyl.}}(u_1, u_2 = u_1^{\dagger}) d\mathbf{u}_1$$

$$= \sum_R \int \chi_R(u_1) \chi_R^{\dagger}(u_1) e^{-A \Omega(R)}$$

$$= \sum_R e^{-A \Omega(R)}$$

Z₃ holeymer: (u₁, u₂, u₃)

Ex:



→ 2-cell with identif

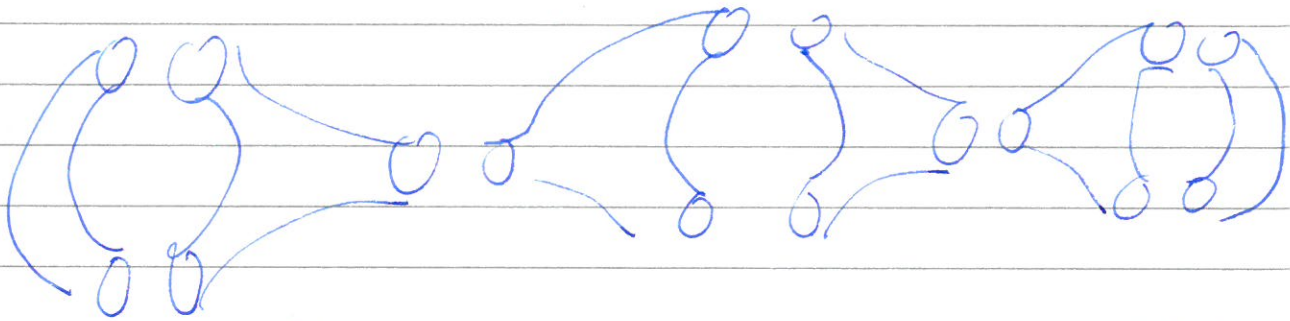
$$\sum_{\mathcal{R}} \int dV e^{-A \rho_2(\mathcal{R})} \chi_{\mathcal{R}}(u_1, u_2, u_3) (v_1, v_2, v_3)$$

Dim \mathcal{R} .

$$\hookrightarrow \sum_{\mathcal{R}} e^{-A G_2(\mathcal{R})} (\text{Dim } \mathcal{R})^{-1} \chi_{\mathcal{R}}(u_1) \chi_{\mathcal{R}}(u_2) \chi_{\mathcal{R}}(u_3)$$

I. 9.

genus 2:



derivation:

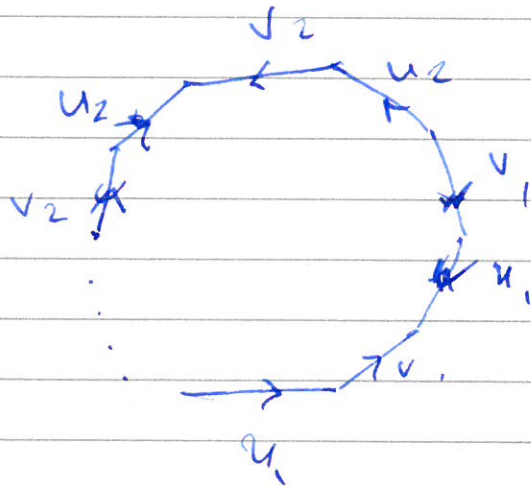
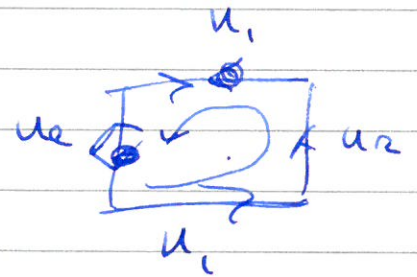
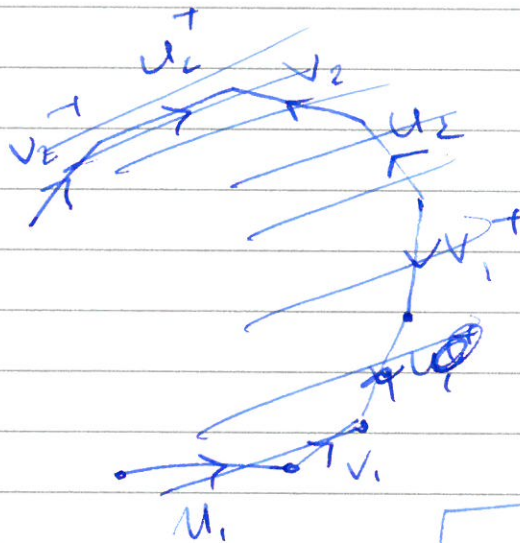
$$Z(G=2, A) = \sum_R (\dim R)^{2-2G} e^{-A Q(R)}$$

For genus G:

$$Z(G, A) = \sum_R (\dim R)^{2-2G} e^{-A Q(R)}$$

Another way:

① A genus g G surface:



$$[u, v] = uvu^+v^+$$

$$\sum_R (\text{Dim } R)^{2g} \chi_R([u_1, v_1] [u_2, v_2] \dots [u_g, v_g])$$

$\rho = \text{Ar}_2(\mathbb{R}) \rightarrow$ same answer.

\mathbb{F} : \mapsto group algebras . (S_n)

\mapsto Centre . $Z(\mathbb{F}(S_n))$

\mapsto \otimes Conj. classes \otimes

. Projectors

\mapsto Schur-Weyl duality

\otimes $(\text{Dim } R)$

\mapsto S_n :

- group $n!$

- all HC-arrangements

1 2 3

↓

1 2 3

(1)(2)(3)

Cycle notation

• $0 \in S_n$:

$\mathbb{Q}(S_n)$: - is an algebra

- A vector space:

- can scale

- can add.

(11.2)

$$a = \sum_{\sigma \in S_n} a_\sigma \sigma$$

$$a \in \mathbb{C}(S_n)$$

$$\text{Dim.} = n!$$

$$\text{Dim} = |G|$$

$$b = \sum_{\sigma \in S_n} b_\sigma \sigma$$

$$\begin{aligned} a \cdot b &= \sum_{\sigma, \tau} a_\sigma b_\tau \underbrace{\sigma \cdot \tau} \\ &= \sum \end{aligned}$$

Associative product.

(H.P.)

$$\begin{aligned} \rightarrow \delta(\sigma) &= 1 \quad \text{if } \sigma = \text{id.} \\ &= 0 \quad \text{if } \sigma \neq \text{id.} \end{aligned}$$

Regular Rep.:

$$|V^{\text{reg.}}| = n!$$

~~reg.~~

$$\mapsto \delta(\sigma) = \frac{1}{n!} \sum_{R} d_R \chi_R(\sigma)$$

- R is an irrep.
- d_R is the dimension of the irrep.

II Pf.

→ Centre:

given G .

$\gamma \in \gamma^{-1}$ is said to be conjugate

(5)

$[(1) (2) (3)]$

$[\begin{array}{l} (12) (3) \\ (23) (1) \\ (12) (3) \end{array}]$

$[\begin{array}{l} (123) \\ (132) \end{array}]$

3. conj. classes

given $\sigma, \tau \in K$.

$\exists \gamma \in G$ s.t. $\gamma \sigma \gamma^{-1} = \tau$

II.5:

$$\rightarrow \underline{\text{Center:}} \text{ of } (R_n) \\ = Z(R_n)$$

⊕ - Vector Subspace

- Any el. of v.s. commutes with all (R_n) .

⊕ Sum over conj. class.

$$\hat{K} = \left(\sum_{\sigma \in K} \sigma \right)$$

⊗

\hat{K} is central element.

$$\oplus \quad \hat{\sigma} = \sum_{\gamma \in R_n} \gamma \sigma \gamma^{-1}$$

is a central element.

$$\text{den } Z(\mathbb{P}(\mathbb{Z}_n))$$

$$= p(n)$$

of partitions of n .

• irreducible characters are
in 1-1 corr. with cong. classes

$$\begin{aligned} \chi^R(0) &= \chi^R(\sigma \sigma^{-1}) \\ &= \chi^R_K \\ &= \end{aligned}$$

$$\# \text{ of } R = \# \text{ of } K$$

$$\# \text{ of irreps of } \mathbb{Z}_n = \cancel{\# \text{ of}} p(n).$$

(11.7)

→ ~~⊗~~ For each irrep —

→ There is a central element

→ in fact —

There is a projector

→ $\mathbb{C}(G)$ spanned by

$$\{ P_{\mathbb{R}} \mid \mathbb{R} \vdash n \}$$

$$P_{\mathbb{R}} = \frac{d_{\mathbb{R}}}{n!} \sum_{\sigma \in S_n} \chi_{\mathbb{R}}(\sigma) \sigma$$

⊗ $P_{\mathbb{R}}^2 = P_{\mathbb{R}}$

$$P_{\mathbb{R}} P_{\mathbb{S}} = \delta_{\mathbb{R}\mathbb{S}} P_{\mathbb{R}}$$

(11.8)

$$\rightarrow \ell = \sum_k c_k k$$

$$= \sum_k \tilde{c}_k P_k$$

$$\ell \sigma = \sigma \ell \quad \forall \sigma.$$

\rightarrow

$$\chi_R(\ell \tau) = \frac{\chi_R(\ell) \chi_R(\tau)}{d_R}$$

(*)

(Ex) Folge aus Schur's Lemma.

\rightarrow

$$\mathbb{D}^R(\ell) = c \mathbb{D}^R(1)$$

(Ex)

Use trick - proof above.

(11.9)

Swt. duality:

$$\rightarrow V_N = \text{Span} \{e_1, e_2, \dots, e_N\}.$$

N -dim. fundamental rep. of $U(N)$

$$\text{Span} \left\{ (e_{i_1} \otimes e_{i_2} \otimes e_{i_3} \otimes \dots \otimes e_{i_n}) \right\}$$

$$= V_N^{\otimes n}$$

\rightarrow How do we decompose this into irreps?

$$\underline{V_N^{\otimes 2}}: \quad \frac{1}{2} (e_{i_1} \otimes e_{i_2} + e_{i_2} \otimes e_{i_1}) \quad \square \square$$

$$\frac{1}{2} (e_{i_1} \otimes e_{i_2} - e_{i_2} \otimes e_{i_1}) \quad \square \cdot$$

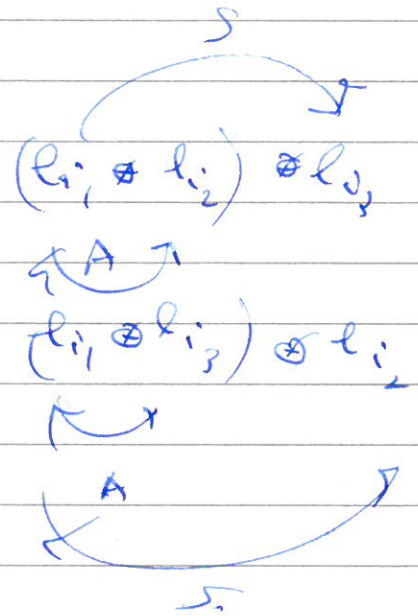
(1)

Linear combinations of def. symm. type

are irreps:

$$\underline{V_H \otimes \rho}$$

has 3 sym. types



2 distinct copies of same rep.

(2) \rightarrow is the dim. of the non-trivial irrep of $\underline{S_3}$

$$3! = 1 + 1 + 4 = \underline{6} = \sum_k d_k^2$$

→ The general connection:

$$V_N^{\otimes \mathbb{R}} = \bigoplus_{\substack{R \\ \ell(R) \leq N}} V_R^{U(R)} \otimes V_R^{F_n}$$

$$\rightarrow \textcircled{a} \quad \underline{U(N)} \quad \& \quad \underline{C(\mathbb{D}_N)}$$

commute with each other
in $V_N^{\otimes \mathbb{R}}$.

~~can be seen~~

① The full Comm of $(\mathbb{C}^N) \rightarrow$
 $U(N)$

② (image of en. algeb of $u(N)$)

③ weights of \underline{h} are labelled by

$\gamma \cdot \lambda$ with \underline{n} boxed

$\mathcal{D} \rightarrow$ be decomp.

\vdash

\mathcal{R}

$\text{Dim } \mathcal{R}$

$d_{\mathcal{R}}$

$$\text{Dim } \mathcal{R} = \frac{\prod_{\substack{\mathcal{N} \\ \mathcal{Q}}} (\mathcal{N} + \mathcal{Q}_0)}{\prod \text{hooks}}$$

$$d_{\mathcal{R}} = \frac{n!}{\prod \text{hooks}}$$

$$\rightarrow \text{tr}_{\mathbb{Z}^{\otimes n}} (1 \otimes P_{\mathbb{R}})$$

$$= \sum_{\sigma} \frac{d_{\mathbb{R}} \chi_{\mathbb{R}}(\sigma)}{n!} \text{tr}_{\mathbb{Z}^{\otimes n}}(\sigma)$$

$$= \frac{d_{\mathbb{R}}}{n!} \sum_{\sigma} \chi_{\mathbb{R}}(\sigma) N_{\sigma}$$

$$= d_{\mathbb{R}} \cdot \text{Dim } \mathbb{R}$$

$$\Rightarrow (\text{Dim } \mathbb{R}) = \left(\frac{N^n}{n!} \right) \chi_{\mathbb{R}}(\Omega)$$

$$\Omega = \sum_{\sigma} \sigma N_{\sigma}$$

$$\Omega \in \mathbb{Z}(\mathbb{C}(S_n))$$

II.14.

$$\rightarrow \text{tr}_{\mathbb{N}} (U \otimes P_R)$$

$$= \text{dr } \chi_R(U)$$

$$= \frac{\text{dr}}{n!} \sum_{\sigma} \chi_R(\sigma) \text{tr} (U^{\otimes n} \otimes \sigma)$$

$$= \frac{\text{dr}}{n!} \sum_{\sigma} \chi_R(\sigma) \prod_i (\text{tr } U^i)^{C_i(\sigma)}$$

$$\chi_R(U) = \frac{1}{n!} \sum_{\sigma} \chi_R(\sigma) \underbrace{P_{\sigma}(e^{\vec{\sigma}})}_{\text{Sym. poly.}}$$

Link: \rightarrow II.13 & II.14 :

should be explained slowly

$$\rightarrow \chi_R(e^{\tau}) = \chi_R(e) \frac{\chi_R(\tau)}{\text{dr}}$$

Exp. slowly ;

\rightarrow Thy. of Sym. Functions ;

1015.

III : Large N expansion:

Z_{char} v/s $Z_{\text{non-char}}$

$$Z_{\text{char}} = \sum_{n=0}^{\infty} \sum_{R \vdash n} (\text{Dim } R)^{2-2G}$$

$$= \sum_{n=0}^{\infty} \sum_{R \vdash n} N^{n(2-2G)} \frac{(\chi_R(\Omega))^{2-2G}}{(n!)^{2-2G}}$$

$G=0$:

$$\frac{N^{2n}}{(n!)^2} (\chi_R(\Omega))^2$$

$$= \frac{N^{2n}}{(n!)^2} \int_{\mathcal{R}} d\mathcal{R} \chi_{\mathcal{R}}(\Omega^2)$$

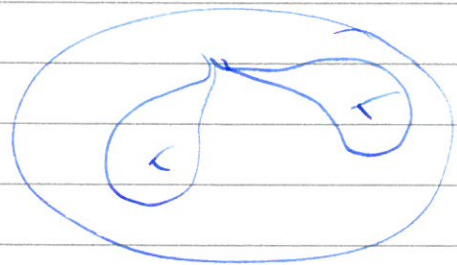
$$Z = \frac{N^{2n}}{(n!)^2} \int_{\mathcal{R}} d\mathcal{R} \chi_{\mathcal{R}}(\Omega^2)$$

$$= \frac{N^{2n}}{(n!)} \delta(\Omega^2)$$

$$= \frac{N^{2_1}}{n!} \sum_{\sigma_1, \sigma_2} \delta(\sigma_1 \sigma_2) N^{\binom{\sigma_1-1}{\sigma_1} + \binom{\sigma_2-1}{\sigma_2}}$$

Counting parts of permutations

⊗ Count branched covers of \mathbb{S}^2
with 2 branch points



⊗ Take \mathbb{S}^2 w/ 2 branch points:

- Label inv. image of gen. pt
get σ_1, σ_2 .

- perm. for whole should be known.

◦ Fermat was:

— Branched covers of S^2

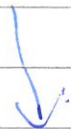
with k branch pts

are counted by

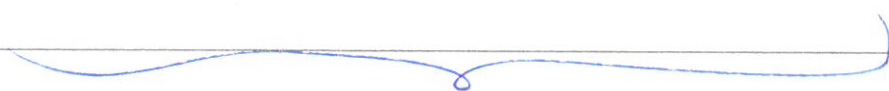
~~counting k -tuples of~~
~~perms.~~

— think about

$$\text{Hom} \left(\pi_1 \left(S^2 \setminus \{k \text{ pts}\} \right) \right)$$



$$S_k$$



→ For generic case:

$$\Omega = \left(\mathbb{1} + \frac{\sigma}{\tau} \right) \quad \text{still}$$

$$\Omega^{-1} = \left(\mathbb{1} + \frac{\sigma}{\tau} \right)^{-1}$$

Expand. (Binomial)

↓ use $\sigma^2 = \tau$.

$$= \frac{1}{\left(1 - \frac{1}{\tau^2}\right)} + \frac{1}{\tau} \frac{\sigma}{\left(1 - \frac{1}{\tau^2}\right)}$$

so

~~$$\Omega^{-1} = \mathbb{1} - \frac{\sigma}{\tau} + \frac{\sigma^2}{\tau^2} - \frac{\sigma^3}{\tau^3} + \dots$$~~

$$\Omega = \sum_{\sigma} N_{\sigma}^{\sigma} \sigma$$

$$= 1 + \sum_{\sigma} N_{\sigma}^{\sigma} \sigma$$

$$\Omega^{-1} = 1 - \sum_{\sigma} N_{\sigma}^{\sigma} \sigma$$

$$+ \sum_{\sigma_1, \sigma_2} N_{\sigma_1}^{\sigma_1} \cdot N_{\sigma_2}^{\sigma_2} \sigma_1 \sigma_2$$

+ ...

$$+ \dots$$

$$\frac{d}{dr} (\Omega \Omega^{-1}) = 1$$

$$\Rightarrow \frac{d\Omega}{dr} \cdot \Omega^{-1} + \Omega \cdot \frac{d\Omega^{-1}}{dr} = 1 \Rightarrow \frac{d\Omega^{-1}}{dr} = -\frac{d\Omega}{dr} \cdot \Omega^{-1}$$

$$\sum_R \frac{Z^{n(2-2G)}}{(n!)^{2-2G}} \left(\chi_R(\Omega) \right)^{2-2G}$$

$$= \sum_R \frac{Z^{n(2-2G)}}{(n!)^{2-2G}} \left(\chi_R(\Omega)^{-1} \right)^{2G-2}$$

$$= \sum_R \frac{Z^{n(2-2G)}}{(n!)^{2-2G}} \left(\frac{\chi_R(\Omega^{-1})}{d_R^2} \right)^{2G-2}$$

$$= \sum_R \frac{Z^{n(2-2G)}}{(n!)^{2-2G}} d_R^{2-2G} \left(\frac{\chi_R(\Omega^{-1})}{d_R} \right)^{2G-2}$$

$$= \sum_R \frac{Z^{n(2-2G)}}{(n!)^{2-2G}} d_R^{2-2G} \left(\frac{\chi_R(\Omega^{-1})}{d_R} \right)^{2G-2}$$

$$\rightarrow \frac{(n!)^2}{d_R^2} = \sum_{s, t} \chi_R(s t s^{-1} t^{-1})$$

$$\text{RHS} = \sum_{s, t} \chi_R(s t s^{-1}) \chi_R(t^{-1})$$

$$= (n!) \sum_{t^{-1}} \frac{\chi_R(t) \chi_R(t^{-1})}{d_R}$$

$$= \frac{(n!)^2}{(d_R^2)}$$

$$\text{LHS} = \sum_{s, t} s t s^{-1} t^{-1}$$

$$\sum_R \left(\frac{\chi_R(t)}{\bar{d}_R} \right)^{2G} \cdot \frac{\chi_R(\Omega^{2-2G})}{\bar{d}_R} \frac{d_R^2}{(n!)^2}$$

$$N^{n(2-2G)}$$

$$= \sum_R \frac{\chi_R(t^G \cdot \Omega^{2-2G})}{\bar{d}_R} \frac{d_R^2}{(n!)^2}$$

$$N^{n(2-2G)}$$

$$= \frac{1}{n!} \delta(t^G \cdot \Omega^{2-2G}) N^{n(2-2G)}$$

$$= \frac{1}{n!} \sum_{\substack{s_1 t_1 \dots \\ s_g t_g}} \delta \left(s_1 t_1 s_1^{-1} t_1^{-1} \dots s_g t_g s_g^{-1} t_g^{-1} \right) \Omega^{2-2g}$$

$$N^{n(2-2g)}$$

①: $\Omega = 1$

- Counting un-branched covers

- $\text{Hom}(\pi_1(\Sigma) \rightarrow S_n)$

- $\frac{N}{n}$

②: When expand: $\chi(F_n)$

$$\frac{1}{n!} \sum_{\substack{s_1 t_1 \dots s_g t_g \\ \sigma_1 \dots \sigma_L}} \delta \left(\pi([s_i, t_i] \sigma_1 \dots \sigma_L) \right) N^{n(2-2g) - 2bc}$$

$d(g, L)$
 $\rightsquigarrow \mathcal{Z}(g, L)$